# Multiple nodal solutions to a Robin problem with sign-changing potential and locally defined reaction 

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#### Abstract

A Robin boundary-value problem with non-homogeneous differential operator, indefinite potential, and reaction defined only near zero is investigated. The existence of one or more nodal solutions is achieved by using truncation, perturbation, and comparison techniques, results from Morse theory, besides variational methods.


Keywords: Robin problem, nodal solutions, indefinite potential, locally defined reaction.

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## 1 Introduction

The main purpose of this paper is to investigate both existence and multiplicity of nodal $C^{1}$-solutions to the following Robin boundary-value problem:

[^0]\[

$$
\begin{cases}-\operatorname{div}(a(\nabla u))+\alpha(x)|u|^{p-2} u=f(x, u) & \text { in } \Omega  \tag{1.1}\\ \frac{\partial u}{\partial n_{a}}+\beta(x)|u|^{p-2} u=0 & \text { on } \partial \Omega\end{cases}
$$
\]

Here, $\Omega$ denotes a bounded domain in $\mathbb{R}^{N}, N \geq 3$, with a smooth boundary $\partial \Omega$, the coefficient $\alpha$ is essentially bounded but sign-changing, $\beta$ lies in $C^{0, \gamma}(\partial \Omega)$ and takes nonnegative values, $1<p<+\infty$, the reaction $f: \Omega \times[-\theta, \theta] \rightarrow \mathbb{R}$ satisfies Carathéodory's conditions. Moreover, $a: \mathbb{R}^{N} \rightarrow \mathbb{R}^{N}$ indicates a strictly monotone map having appropriate regularity and growth properties, while $\frac{\partial}{\partial n_{a}}$ stands for the co-normal derivative associated with $a$; cf. Section 2.
Problem (1.1) exhibits at least three interesting features:
i) We do not require that $\xi \mapsto a(\xi)$ be ( $p-1$ )-homogeneous. So, meaningful differential operators, as the $(p, q)$-Laplacian, are incorporated in (1.1).
ii) The potential term $\alpha(x)|u|^{p-2} u$ turns out indefinite, because $\alpha(x)$ can change sign.
iii) $t \mapsto f(x, t)$ is only locally defined, whence its behavior near zero matters, and no conditions at infinity are imposed.
Via truncation-perturbation-comparison techniques, results from Morse theory, besides variational methods, we obtain a nodal solution $\hat{u} \in C^{1}(\bar{\Omega})$ of (1.1); see Theorems 3.13.2 below. The case $p>2$ and $a(\xi):=\left(|\xi|^{p-2}+1\right) \xi$, namely when the $(p, 2)$-Laplacian appears, is examined next in Theorem 3.3, which allows also $f$ to be resonant.

As far as we know, the existence of sign-changing solutions to Robin problems that exhibit difficulties i)-iii) did not receive much attention up to now. Topics i) and, somehow, iii) have been recently addressed in [25], while [2, 18] investigate ii) but for $a(\xi):=|\xi|^{p-2} \xi$. Further items can evidently be found in their bibliographies.

Section 4 deals with multiplicity. If $f(x, \cdot)$ is odd, Theorem 4.1 gives a whole sequence $\left\{u_{n}\right\} \subseteq C^{1}(\bar{\Omega})$ of nodal solutions to (1.1) such that $u_{n} \rightarrow 0$ in $C^{1}(\bar{\Omega})$. The works $[9,14]$ contain similar results concerning Dirichlet problems without indefinite potential. All of them exploit an abstract theorem by Kajikiya [11]. Once the map $a$ is particularized, we can do without symmetry and still produce two or three nodal $C^{1}$-solutions; cf. Theorem 4.2 , which treats ( $p, 2$ )-Laplace equations, and Theorem 4.3, where in addition $p=2$. It should be pointed out that one solution always comes from a flow invariance argument patterned after that of [8], devoted to Neumann's case; see also [21].

## 2 Preliminaries. The map $\xi \mapsto a(\xi)$

Let $(X,\|\cdot\|)$ be a real Banach space. Given a set $V \subseteq X$, write $\bar{V}$ for the closure of $V$, $\partial V$ for the boundary of $V$, and $\operatorname{int}_{X}(V)$ or simply $\operatorname{int}(V)$, when no confusion can arise,
for the interior of $V$. If $x \in X$ and $\delta>0$ then

$$
B_{\delta}(x):=\{z \in X:\|z-x\|<\delta\}, \quad B_{\delta}:=B_{\delta}(0)
$$

The symbol $\left(X^{*},\|\cdot\|_{X^{*}}\right)$ denotes the dual space of $X,\langle\cdot, \cdot\rangle$ indicates the duality pairing between $X$ and $X^{*}$, while $x_{n} \rightarrow x$ (respectively, $x_{n} \rightharpoonup x$ ) in $X$ means 'the sequence $\left\{x_{n}\right\}$ converges strongly (respectively, weakly) in $X^{\prime}$. We say that $A: X \rightarrow X^{*}$ is of type ( S$)_{+}$ provided

$$
x_{n} \rightharpoonup x \text { in } X, \quad \limsup _{n \rightarrow+\infty}\left\langle A\left(x_{n}\right), x_{n}-x\right\rangle \leq 0 \quad \Longrightarrow \quad x_{n} \rightarrow x .
$$

The function $\Phi: X \rightarrow \mathbb{R}$ is called coercive if $\lim _{\|x\| \rightarrow+\infty} \Phi(x)=+\infty$ and weakly sequentially lower semi-continuous when

$$
x_{n} \rightharpoonup x \text { in } X \quad \Longrightarrow \quad \Phi(x) \leq \liminf _{n \rightarrow \infty} \Phi\left(x_{n}\right) .
$$

Let $\Phi \in C^{1}(X)$. The classical Palais-Smale compactness condition for $\Phi$ reads as follows.
(PS) Every sequence $\left\{x_{n}\right\} \subseteq X$ such that $\left\{\Phi\left(x_{n}\right)\right\}$ is bounded and $\lim _{n \rightarrow+\infty}\left\|\Phi^{\prime}\left(x_{n}\right)\right\|_{X^{*}}=0$ has a convergent subsequence.

The next elementary result [15, Proposition 2.2] will be employed later.
Proposition 2.1. Suppose $X$ reflexive, $\Phi \in C^{1}(X)$ coercive, and $\Phi^{\prime}=A+B$, with $A$ of type $(\mathrm{S})_{+}$and $B$ compact. Then $\Phi$ satisfies (PS).

Define, for every $c \in \mathbb{R}$,

$$
\Phi^{c}:=\{x \in X: \Phi(x) \leq c\}, \quad K_{c}(\Phi):=K(\Phi) \cap\{x \in X: \Phi(x)=c\}
$$

where, as usual, $K(\Phi)$ denotes the critical set of $\Phi$, i.e., $K(\Phi):=\left\{x \in X: \Phi^{\prime}(x)=0\right\}$.
Given a topological pair $(A, B)$ fulfilling $B \subset A \subseteq X$, the symbol $H_{k}(A, B), k \in \mathbb{N}_{0}$, indicates the $\mathrm{k}^{\text {th }}$-relative singular homology group of $(A, B)$ with integer coefficients. If $x_{0} \in K_{c}(\Phi)$ is an isolated point of $K(\Phi)$ then

$$
C_{k}\left(\Phi, x_{0}\right):=H_{k}\left(\Phi^{c} \cap V, \Phi^{c} \cap V \backslash\left\{x_{0}\right\}\right), \quad k \in \mathbb{N}_{0}
$$

are the critical groups of $\Phi$ at $x_{0}$. Here, $V$ stands for any neighborhood of $x_{0}$ such that $K(\Phi) \cap \Phi^{c} \cap V=\left\{x_{0}\right\}$. By excision, this definition does not depend on the choice of $V$. Suppose $\Phi$ satisfies condition (PS), $\left.\Phi\right|_{K(\Phi)}$ is bounded below, and $c<\inf _{x \in K(\Phi)} \Phi(x)$. Put

$$
C_{k}(\Phi, \infty):=H_{k}\left(X, \Phi^{c}\right), \quad k \in \mathbb{N}_{0} .
$$

The second deformation lemma [4, Theorem 5.1.33] implies that this definition does not depend on the choice of $c$. If $K(\Phi)$ is finite, then setting

$$
M(t, x):=\sum_{k=0}^{+\infty} \operatorname{rank} C_{k}(\Phi, x) t^{k}, \quad P(t, \infty):=\sum_{k=0}^{+\infty} \operatorname{rank} C_{k}(\Phi, \infty) t^{k} \quad \forall(t, x) \in \mathbb{R} \times K(\Phi),
$$

the following Morse relation holds:

$$
\begin{equation*}
\sum_{x \in K(\Phi)} M(t, x)=P(t, \infty)+(1+t) Q(t), \tag{2.1}
\end{equation*}
$$

where $Q(t)$ denotes a formal series with nonnegative integer coefficients; see for instance [19, Theorem 6.62].

Now, let $X$ be a Hilbert space, let $x \in K(\Phi)$, and let $\Phi$ be $C^{2}$ in a neighborhood of $x$. If $\Phi^{\prime \prime}(x)$ turns out to be invertible, then $x$ is called non-degenerate. The Morse index $d$ of $x$ is the supremum of the dimensions of the vector subspaces of $X$ on which $\Phi^{\prime \prime}(x)$ turns out to be negative definite. When $x$ is non-degenerate and with Morse index $d$ one has $C_{k}(\Phi, x)=\delta_{k, d} \mathbb{Z}, k \in \mathbb{N}_{0}$. The monograph [19] represents a general reference on the subject.

Henceforth, $\Omega$ will denote a bounded domain of the real Euclidean space ( $\mathbb{R}^{N},|\cdot|$ ), $N \geq 3$, with a $C^{2}$-boundary $\partial \Omega$, on which we will employ the $(N-1)$-dimensional Hausdorff measure $\sigma$. The symbol $n(x)$ indicates the outward unit normal vector to $\partial \Omega$ at its point $x, p \in] 1,+\infty\left[, p^{\prime}:=p /(p-1),\|\cdot\|_{s}\right.$ with $s \geq 1$ is the usual norm of $L^{s}(\Omega)$, $X:=W^{1, p}(\Omega)$, and

$$
\|u\|:=\left(\|u\|_{p}^{p}+\|\nabla u\|_{p}^{p}\right)^{1 / p}, \quad u \in X .
$$

Write $p^{*}$ for the critical exponent of the Sobolev embedding $W^{1, p}(\Omega) \hookrightarrow L^{s}(\Omega)$. Recall that $p^{*}=N p /(N-p)$ if $p<N, p^{*}=+\infty$ otherwise, and the embedding turns out to be compact whenever $1 \leq s<p^{*}$. Given $t \in \mathbb{R}$ and $u, v: \Omega \rightarrow \mathbb{R}, t^{ \pm}:=\max \{ \pm t, 0\}$, $u^{ \pm}(x):=u(x)^{ \pm}, u \leq v$ (resp., $u<v$, etc.) means $u(x) \leq v(x)$ (resp., $u(x)<v(x)$, etc.) for almost every $x \in \Omega$. If $u, v$ belongs to a function space, say $Y$, then we set

$$
[u, v]:=\{w \in Y: u \leq w \leq v\}, \quad Y_{+}:=\{w \in Y: w \geq 0\} .
$$

Putting $C_{+}:=C^{1}(\bar{\Omega})_{+}$and $\operatorname{int}\left(C_{+}\right):=\operatorname{int}_{C^{1}(\bar{\Omega})}\left(C_{+}\right)$, one evidently has

$$
\operatorname{int}\left(C_{+}\right):=\left\{u \in C_{+}: u(x)>0 \forall x \in \bar{\Omega}\right\} .
$$

From now on, $c_{1}, c_{2}, \ldots$ denote appropriate positive constants while $\left.\mathbb{R}^{+}:=\right] 0,+\infty[$. Let $\omega \in C^{1}\left(\mathbb{R}^{+}\right)$and let $\tau \in[1, p[$ satisfy

$$
\begin{equation*}
c_{1} \leq \frac{t \omega^{\prime}(t)}{\omega(t)} \leq c_{2}, \quad c_{3} t^{p-1} \leq \omega(t) \leq c_{4}\left(t^{\tau-1}+t^{p-1}\right), \quad t \in \mathbb{R}^{+} . \tag{2.2}
\end{equation*}
$$

The following assumptions on $a: \mathbb{R}^{N} \rightarrow \mathbb{R}^{N}$ will be posited.
$\left(\mathrm{a}_{1}\right) a(\xi):=a_{0}(|\xi|) \xi$ for all $\xi \in \mathbb{R}^{N}$, where $a_{0} \in C^{1}\left(\mathbb{R}^{+}, \mathbb{R}^{+}\right), t \mapsto t a_{0}(t)$ turns out to be strictly increasing on $\mathbb{R}^{+}$, and

$$
\lim _{t \rightarrow 0^{+}} t a_{0}(t)=0, \quad \lim _{t \rightarrow 0^{+}} \frac{t a_{0}^{\prime}(t)}{a_{0}(t)}>-1 .
$$

$\left(\mathrm{a}_{2}\right)|\nabla a(\xi)| \leq c_{5} \frac{\omega(|\xi|)}{|\xi|}$ for every $\xi \in \mathbb{R}^{N} \backslash\{0\}$.
$\left(\mathrm{a}_{3}\right)(\nabla a(\xi) y) \cdot y \geq \frac{\omega(|\xi|)}{|\xi|}|y|^{2}$ for all $\xi, y \in \mathbb{R}^{N}$ with $\xi \neq 0$.
$\left(\mathrm{a}_{4}\right)$ If $G_{0}(t):=\int_{0}^{t} s a_{0}(s) d s$, then there exist $1<\hat{q}<q \leq p$ such that $\tau \leq q$,

$$
c_{6} t^{p} \leq t^{2} a_{0}(t)-\hat{q} G_{0}(t) \quad \forall t \in \mathbb{R}^{+}, \quad \lim _{t \rightarrow 0^{+}} \frac{q G_{0}(t)}{t^{q}}=c_{7},
$$

and $t \mapsto G_{0}\left(t^{1 / q}\right)$ is convex on $[0,+\infty[$.
Remark 2.1. Conditions $\left(\mathrm{a}_{1}\right)-\left(\mathrm{a}_{3}\right)$ come from Lieberman's nonlinear regularity theory [12] and Pucci-Serrin's maximum principle [26].

Lemma 2.1. Let $\left(\mathrm{a}_{1}\right)-\left(\mathrm{a}_{3}\right)$ be satisfied. Then:
$\left(\mathrm{i}_{1}\right) \xi \mapsto a(\xi)$ is strictly monotone, continuous, and, a fortiori, maximal monotone.
$\left(\mathrm{i}_{2}\right)|a(\xi)| \leq c_{8}\left(|\xi|^{\tau-1}+|\xi|^{p-1}\right)$ for every $\xi \in \mathbb{R}^{N}$.
$\left(\mathrm{i}_{3}\right) a(\xi) \cdot \xi \geq \frac{c_{3}}{p-1}|\xi|^{p}$ for all $\xi \in \mathbb{R}^{N}$.
Proof. Conclusions ( $\mathrm{i}_{1}$ ) $-\left(\mathrm{i}_{2}\right)$ are obvious. Let us verify ( $\mathrm{i}_{3}$ ). Because of ( $\mathrm{a}_{3}$ ) and (2.2) we easily obtain

$$
a(\xi) \cdot \xi=\int_{0}^{1} \frac{d}{d t} a(t \xi) \cdot \xi d t=\int_{0}^{1}(\nabla a(t \xi) \xi) \cdot \xi d t \geq \int_{0}^{1} \frac{\omega(t|\xi|)}{t|\xi|}|\xi|^{2} d t \geq \frac{c_{3}}{p-1}|\xi|^{p}
$$

for every $\xi \in \mathbb{R}^{N}$.
Remark 2.2. Thanks to $\left(\mathrm{a}_{1}\right)$, the function $G_{0}$ defined in $\left(\mathrm{a}_{4}\right)$ is strictly increasing and convex. Consequently, also the map $G: \mathbb{R}^{N} \rightarrow \mathbb{R}$ given by

$$
G(\xi):=G_{0}(|\xi|), \quad \xi \in \mathbb{R}^{N},
$$

turns out convex. Moreover, $G(0)=0$,

$$
\begin{equation*}
\nabla G(\xi)=G_{0}^{\prime}(|\xi|) \frac{\xi}{|\xi|}=a_{0}(|\xi|) \xi=a(\xi) \text { provided } \xi \neq 0 \tag{2.3}
\end{equation*}
$$

as well as, on account of $\left(a_{1}\right)$,

$$
\begin{equation*}
G(\xi) \leq a(\xi) \cdot \xi \quad \forall \xi \in \mathbb{R}^{N} \tag{2.4}
\end{equation*}
$$

Lemma 2.2. If $\left(\mathrm{a}_{1}\right)-\left(\mathrm{a}_{4}\right)$ hold true, then

$$
\begin{equation*}
\frac{c_{3}}{p(p-1)}|\xi|^{p} \leq G(\xi) \leq c_{9}\left[|\xi|^{q}+|\xi|^{p}\right], \quad \xi \in \mathbb{R}^{N} . \tag{2.5}
\end{equation*}
$$

Proof. Through (2.3) and ( $\mathrm{i}_{3}$ ) in Lemma 2.1 we get

$$
G(\xi)=\int_{0}^{1} \frac{d}{d t} G(t \xi) d t=\int_{0}^{1} a(t \xi) \cdot \xi d t \geq \int_{0}^{1} \frac{c_{3}}{p-1}|\xi|^{p} t^{p-1} d t=\frac{c_{3}}{p(p-1)}|\xi|^{p} .
$$

The other inequality easily follows from $\left(\mathrm{a}_{4}\right),(2.4)$, and ( $\mathrm{i}_{2}$ ) of Lemma 2.1.
Example 2.1. The functions $a_{0}$ listed below comply with $\left(\mathrm{a}_{1}\right)-\left(\mathrm{a}_{4}\right)$.

- $a_{0}(t):=t^{p-2}$ for every $t \in \mathbb{R}^{+}$. It corresponds to the well-known $p$-Laplacian $\Delta_{p}$, defined by

$$
\Delta_{p} u:=\operatorname{div}\left(|\nabla u|^{p-2} \nabla u\right), \quad u \in X
$$

- $a_{0}(t):=t^{p-2}+t^{q-2}$ for all $t \in \mathbb{R}^{+}$. The associated operator, usually called $(p, q)-$ Laplacian, arises in mathematical physics; see, e.g., the survey paper [13].
- $a_{0}(t):=\left(1+t^{2}\right)^{\frac{p-2}{2}}$ for every $t \in \mathbb{R}^{+}$. This function stems form the generalized $p$-mean curvature operator, namely

$$
u \mapsto \operatorname{div}\left[\left(1+|\nabla u|^{2}\right)^{\frac{p-2}{2}} \nabla u\right], u \in X .
$$

- $a_{0}(t):=t^{p-2}\left(1+\frac{1}{1+t^{p}}\right)$ for all $t \in \mathbb{R}^{+}$. It corresponds to the differential operator

$$
u \mapsto \Delta_{p} u+\operatorname{div}\left(\frac{|\nabla u|^{p-2}}{1+|\nabla u|^{p}} \nabla u\right), \quad u \in X,
$$

which employs in plasticity theory [3].

Finally, let $A: X \rightarrow X^{*}$ be the nonlinear operator associated with $a$, i.e.,

$$
\langle A(u), v\rangle:=\int_{\Omega} a(\nabla u) \cdot \nabla v d x \quad \forall u, v \in X
$$

Proposition 3.5 in [5] ensures that $A$ is bounded, continuous, monotone, as well as of type $(S)_{+}$. Moreover, via the nonlinear Green's identity [4, Theorem 2.4.54] we easily have Remark 2.3. If $u \in X, w \in L^{p^{\prime}}(\Omega)$, and $\beta \in C^{0, \gamma}\left(\partial \Omega, \mathbb{R}_{0}^{+}\right)$then

$$
\langle A(u), v\rangle+\int_{\partial \Omega} \beta(x)|u(x)|^{p-2} u(x) v(x) d \sigma=\int_{\Omega} w(x) v(x) d x, \quad v \in X,
$$

is equivalent to

$$
-\operatorname{div}(a(\nabla u))=w(x) \quad \text { in } \quad \Omega, \quad \frac{\partial u}{\partial n_{a}}+\beta(x)|u|^{p-2} u=0 \quad \text { on } \quad \partial \Omega .
$$

Here, $\frac{\partial u}{\partial n_{a}}$ denotes the co-normal derivative of $u$, defined extending the map $v \mapsto a(\nabla v) \cdot n$ from $C^{1}(\bar{\Omega})$ to $X$.

Let $\beta \in C^{0, \gamma}\left(\partial \Omega, \mathbb{R}_{0}^{+}\right)$, let $h: \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ be a Carathéodory function such that

$$
|h(x, t)| \leq C\left(1+|t|^{s-1}\right) \quad \forall(x, t) \in \Omega \times \mathbb{R},
$$

where $C>0,1 \leq s \leq p^{*}$, and let $H(x, t):=\int_{0}^{t} h(x, \tau) d \tau$. Consider the $C^{1}$-functional $\varphi_{h}: X \rightarrow \mathbb{R}$ defined by

$$
\varphi_{h}(u):=\int_{\Omega} G(\nabla u(x)) d x+\frac{1}{p} \int_{\partial \Omega} \beta(x)|u(x)|^{p} d \sigma-\int_{\Omega} H(x, u(x)) d x, \quad u \in X .
$$

A relation between local minimizers of $\varphi_{h}$ in $C^{1}(\bar{\Omega})$ and in $X$ occurs [23, Proposition 8].
Proposition 2.2. Suppose $u_{0} \in X$ is a local $C^{1}(\bar{\Omega})$-minimizer to $\varphi_{h}$. Then $u_{0} \in C^{1, \eta}(\bar{\Omega})$ and is a local minimizer of $\varphi_{h}$.

We shall employ some facts about the spectrum of the operator

$$
u \mapsto-\Delta_{q} u+\alpha(x)|u|^{q-2} u
$$

in $X$ with homogeneous Robin boundary conditions. So, consider the eigenvalue problem

$$
\begin{equation*}
-\Delta_{q} u+\alpha(x)|u|^{q-2} u=\lambda|u|^{q-2} u \quad \text { in } \quad \Omega, \quad \frac{\partial u}{\partial n_{q}}+\beta(x)|u|^{q-2} u=0 \quad \text { on } \quad \partial \Omega, \tag{2.6}
\end{equation*}
$$

where, henceforth,

$$
\begin{equation*}
\left.\alpha \in L^{\infty}(\Omega) \text { and } \beta \in C^{0, \gamma}\left(\partial \Omega, \mathbb{R}_{0}^{+}\right) \text {with } \gamma \in\right] 0,1[. \tag{2.7}
\end{equation*}
$$

Define

$$
\begin{equation*}
\mathcal{E}_{q}(u):=\|\nabla u\|_{q}^{q}+\int_{\Omega} \alpha(x)|u(x)|^{q} d x+\int_{\partial \Omega} \beta(x)|u(x)|^{q} d \sigma \quad \forall u \in X . \tag{2.8}
\end{equation*}
$$

The Liusternik-Schnirelman theory provides a strictly increasing sequence $\left\{\hat{\lambda}_{n}(q, \alpha, \beta)\right\}$ of eigenvalues for (2.6). As in [20, 22], one has
$\left(\mathrm{p}_{1}\right) \hat{\lambda}_{1}(q, \alpha, \beta)$ turns out to be isolated and simple. Further, $\hat{\lambda}_{1}(q, \alpha, \beta)=\inf _{u \in X \backslash\{0\}} \frac{\mathcal{E}_{q}(u)}{\|u\|_{q}^{q}}$.
$\left(\mathrm{p}_{2}\right)$ There is an eigenfunction $\hat{u}_{1}(q, \alpha, \beta) \in \operatorname{int}\left(C_{+}\right)$associated with $\hat{\lambda}_{1}(q, \alpha, \beta)$ such that $\left\|u_{1}(q, \alpha, \beta)\right\|_{q}=1$.
$\left(\mathrm{p}_{3}\right)$ Write $U:=\left\{u \in X:\|u\|_{q}=1\right\}$ and

$$
\hat{\Gamma}:=\left\{\hat{\gamma} \in C^{0}([-1,1], U): \hat{\gamma}(-1)=-\hat{u}_{1}(q, \alpha, \beta), \hat{\gamma}(1)=\hat{u}_{1}(q, \alpha, \beta)\right\} .
$$

Then $\hat{\lambda}_{2}(q, \alpha, \beta)=\inf _{\hat{\gamma} \in \hat{\Gamma}} \max _{t \in[-1,1]} \mathcal{E}_{q}(\hat{\gamma}(t))$.
Evidently, the set $U_{C}:=\left\{u \in C^{1}(\bar{\Omega}):\|u\|_{q}=1\right\}$ turns out dense in $U$. Moreover, if

$$
\hat{\Gamma}_{C}:=\left\{\hat{\gamma} \in C^{0}\left([-1,1], U_{C}\right): \hat{\gamma}(-1)=-\hat{u}_{1}(q, \alpha, \beta), \hat{\gamma}(1)=\hat{u}_{1}(q, \alpha, \beta)\right\} .
$$

then (cf. [17, Lemma 2.1])
Lemma 2.3. $\hat{\Gamma}_{C}$ is dense in $\hat{\Gamma}$ with respect to the usual norm of $C^{0}([-1,1], X)$.
Let $q:=2$. Denote by $E\left(\hat{\lambda}_{n}\right)$ the eigenspace coming from $\hat{\lambda}_{n}:=\hat{\lambda}_{n}(2, \alpha, \beta)$. It is known $[1,16]$ that $H^{1}(\Omega)=\oplus_{n=1}^{\infty} E\left(\hat{\lambda}_{n}\right)$ and that
$\left(\mathrm{p}_{4}\right)$ For every $n \geq 2$ one has

$$
\hat{\lambda}_{n}=\sup \left\{\frac{\mathcal{E}_{2}(u)}{\|u\|_{2}^{2}}: u \in \bar{H}_{n}, u \neq 0\right\}=\inf \left\{\frac{\mathcal{E}_{2}(u)}{\|u\|_{2}^{2}}: u \in \hat{H}_{n}, u \neq 0\right\},
$$

where

$$
\bar{H}_{n}:=\oplus_{i=1}^{n} E\left(\hat{\lambda}_{i}\right), \quad \hat{H}_{n}:=\oplus_{i=n}^{\infty} E\left(\hat{\lambda}_{i}\right) .
$$

## 3 Nodal solutions: existence

To avoid unnecessary technicalities, 'for every $x \in \Omega$ ' will take the place of 'for almost every $x \in \Omega$ ' while $C_{1}, C_{2}, \ldots$ indicate positive constants arising from the context.

Let $\underline{u}, \bar{u} \in C^{0}(\bar{\Omega}) \cap W^{1, p}(\Omega)$ satisfy

$$
\max _{x \in \bar{\Omega}} \underline{u}(x)<0<\min _{x \in \bar{\Omega}} \bar{u}(x),
$$

let $\theta:=\max \left\{\|\underline{u}\|_{\infty},\|\bar{u}\|_{\infty}\right\}$, and let $f: \Omega \times[-\theta, \theta] \rightarrow \mathbb{R}$ be a Carathéodory function. We shall make the following two hypotheses throughout the paper.
$\left(f_{1}\right)$ One has

$$
\left\{\begin{array}{l}
\langle A(\underline{u}), v\rangle+\int_{\Omega} \alpha|\underline{u}|^{p-2} \underline{u} v d x \leq 0 \leq \int_{\Omega} f(x, \underline{u}) v d x, \\
\int_{\Omega} f(x, \bar{u}) v d x \leq 0 \leq\langle A(\bar{u}), v\rangle+\int_{\Omega} \alpha|\bar{u}|^{p-2} \bar{u} v d x
\end{array} \quad \forall v \in X_{+} .\right.
$$

Alternatively,

$$
\left\{\begin{array}{l}
\langle A(\underline{u}), v\rangle \leq 0 \leq \int_{\Omega}\left[f(x, \underline{u})-\alpha|\underline{u}|^{p-2} \underline{u}\right] v d x \\
\int_{\Omega}\left[f(x, \bar{u})-\alpha|\bar{u}|^{p-2} \bar{u}\right] v d x \leq 0 \leq\langle A(\bar{u}), v\rangle
\end{array} \quad \forall v \in X_{+} .\right.
$$

$\left(\mathrm{f}_{2}\right)$ There exists $a_{\theta} \in L^{\infty}(\Omega)$ such that $|f(x, t)| \leq a_{\theta}(x)$ in $\Omega \times[-\theta, \theta]$.
Remark 3.1. Condition ( $\mathrm{f}_{1}$ ) evidently entails either $f(\cdot, \bar{u}) \leq 0 \leq f(\cdot, \underline{u})$ or

$$
f(\cdot, \bar{u})-\alpha|\bar{u}|^{p-2} \bar{u} \leq 0 \leq f(\cdot, \underline{u})-\alpha|\underline{u}|^{p-2} \underline{u} .
$$

Different behaviors of $f$ at zero will instead be investigated to get nodal solutions of (1.1).

### 3.1 The ( $q-1$ )-sub-linear case

For $q, \hat{q}$ as in $\left(\mathrm{a}_{4}\right)$ and uniformly with respect to $x \in \Omega$, assume that:
( $\mathrm{f}_{3}$ ) $\lim _{t \rightarrow 0} \frac{f(x, t)}{|t|^{q-2} t}=+\infty$.
$\left(\mathrm{f}_{4}\right) \lim _{t \rightarrow 0} \frac{\hat{q} F(x, t)-f(x, t) t}{|t|^{p}}>\left(1-\frac{\hat{q}}{p}\right)\|\alpha\|_{\infty}$.

Example 3.1. Let $\alpha \neq 0$. If $\underline{u}(x):=-\|\alpha\|_{\infty}, \bar{u}(x):=\|\alpha\|_{\infty}$, and

$$
f(x, t):=C_{1}|t|^{q-2-\varepsilon} t-C_{2}|t|^{q-2} t, \quad(x, t) \in \Omega \times[-\theta, \theta]
$$

where $\varepsilon>0$ is small while $C_{1}, C_{2}>0$ comply with $f(\cdot, \bar{u}) \leq-\|\alpha\|_{\infty}^{p} \leq f(\cdot, \underline{u})$, then $\left(\mathrm{f}_{1}\right)-\left(\mathrm{f}_{4}\right)$ hold true.
Remark 3.2. Because of $\left(\mathrm{f}_{2}\right)-\left(\mathrm{f}_{3}\right)$, given $r>p$, to each $\eta>0$ there corresponds $C_{\eta}>0$ fulfilling

$$
\begin{equation*}
f(x, t) t \geq \eta|t|^{q}-C_{\eta}|t|^{r} \quad \forall(x, t) \in \Omega \times[-\theta, \theta] . \tag{3.1}
\end{equation*}
$$

Define, provided $(x, t) \in \Omega \times \mathbb{R}$,

$$
h(x, t):= \begin{cases}\left(\eta|\underline{u}(x)|^{q-2}-C_{\eta}|\underline{u}(x)|^{r-2}+|\underline{u}(x)|^{p-2}\right) \underline{u}(x) & \text { when } t<\underline{u}(x),  \tag{3.2}\\ \eta|t|^{q-2} t-C_{\eta}|t|^{r-2} t+|t|^{p-2} t & \text { if } \underline{u}(x) \leq t \leq \bar{u}(x), \\ \eta \bar{u}(x)^{q-1}-C_{\eta} \bar{u}(x)^{r-1}+\bar{u}(x)^{p-1} & \text { when } t>\bar{u}(x),\end{cases}
$$

and $H(x, t):=\int_{0}^{t} h(x, s) d s$, besides

$$
b(x, t):= \begin{cases}\beta(x)|\underline{u}(x)|^{p-2} \underline{u}(x) & \text { for } t<\underline{u}(x)  \tag{3.3}\\ \beta(x)|t|^{p-2} t & \text { if } \underline{u}(x) \leq t \leq \bar{u}(x) \\ \beta(x) \bar{u}(x)^{p-1} & \text { for } t>\bar{u}(x)\end{cases}
$$

and $B(x, t):=\int_{0}^{t} b(x, s) d s$, where $(x, t) \in \partial \Omega \times \mathbb{R}$. Consider the auxiliary Robin problem

$$
\left\{\begin{array}{l}
-\operatorname{div}(a(\nabla u))+(|\alpha(x)|+1)|u|^{p-2} u=h(x, u) \text { in } \Omega  \tag{3.4}\\
\frac{\partial u}{\partial n_{a}}+b(x, u)=0 \text { on } \partial \Omega
\end{array}\right.
$$

Lemma 3.1. Let $\left(\mathrm{a}_{1}\right)-\left(\mathrm{a}_{4}\right)$, (2.7), $\left(\mathrm{f}_{1}\right)-\left(\mathrm{f}_{3}\right)$ be satisfied. Then (3.4) possesses a unique positive solution $u_{*} \in[0, \bar{u}] \cap \operatorname{int}\left(C_{+}\right)$and a unique negative solution $v_{*} \in[\underline{u}, 0] \cap\left(-\operatorname{int}\left(C_{+}\right)\right)$. Proof. Thanks to (2.5), (3.2), and (3.3), the $C^{1}$-functional $\psi_{h}: X \rightarrow \mathbb{R}$ defined by

$$
\psi_{h}(u):=\int_{\Omega} G(\nabla u) d x+\frac{1}{p} \int_{\Omega}(|\alpha|+1)|u|^{p} d x+\int_{\partial \Omega} B(x, u) d \sigma-\int_{\Omega} H\left(x, u^{+}\right) d x, \quad u \in X,
$$

is coercive. A standard argument, which exploits Sobolev's embedding theorem and the compactness of the trace operator, ensures that $\psi_{h}$ is weakly sequentially lower semicontinuous. Hence,

$$
\begin{equation*}
\inf _{u \in X} \psi_{h}(u)=\psi_{h}\left(u_{*}\right) \tag{3.5}
\end{equation*}
$$

for some $u_{*} \in X$. One has $u_{*} \neq 0$. Indeed, ( $\mathrm{a}_{4}$ ) yields $\delta>0$ fulfilling

$$
\left.\left.G_{0}(t)<\left(c_{7}+1\right) \frac{t^{q}}{q} \quad \forall t \in\right] 0, \delta\right]
$$

Obviously, we may suppose $\delta \leq \min \left\{1, \min _{x \in \bar{\Omega}} \bar{u}(x)\right\}$. If $\rho>0$ is so small that

$$
0<\rho \hat{u}_{1}(x) \leq \delta, \quad \rho\left|\nabla \hat{u}_{1}(x)\right| \leq \delta, \quad x \in \bar{\Omega}
$$

where $\hat{u}_{1}$ comes from $\left(\mathrm{p}_{2}\right)$ with $\alpha_{0}:=|\alpha| / c_{7}$ and $\beta_{0}:=\beta / c_{7}$ in place of $\alpha$ and $\beta$ (vide Section 2), respectively, then

$$
\begin{aligned}
\psi_{h}\left(\rho \hat{u}_{1}\right) & \leq\left(c_{7}+1\right) \frac{\rho^{q}}{q}\left[\left\|\nabla \hat{u}_{1}\right\|_{q}^{q}+\int_{\Omega} \alpha_{0} \hat{u}_{1}^{q} d x+\int_{\partial \Omega} \beta_{0} \hat{u}_{1}^{q} d \sigma\right]-\frac{\eta}{q} \rho^{q}+\frac{C_{\eta}}{r}\left\|\hat{u}_{1}\right\|_{r}^{r} \rho^{r} \\
& =\frac{\rho^{q}}{q}\left[\left(c_{7}+1\right) \hat{\lambda}_{1}\left(q, \alpha_{0}, \beta_{0}\right)-\eta\right]+\frac{C_{\eta}}{r}\left\|\hat{u}_{1}\right\|_{r}^{r} \rho^{r},
\end{aligned}
$$

because $q \leq p, \delta \leq 1,\left\|\hat{u}_{1}\right\|_{q}=1$. Choosing $\eta>\left(c_{7}+1\right) \hat{\lambda}_{1}\left(q, \alpha_{0}, \beta_{0}\right)$ in (3.1), recalling that $r>q$, and decreasing $\rho$ when necessary, entails $\psi_{h}\left(\rho \hat{u}_{1}\right)<0$, whence $u_{*} \neq 0$, as an easy contradiction argument shows.

Through (3.5) we next get $\psi_{h}^{\prime}\left(u_{*}\right)=0$, namely

$$
\begin{equation*}
\left\langle A\left(u_{*}\right), w\right\rangle+\int_{\Omega}(|\alpha|+1)\left|u_{*}\right|^{p-2} u_{*} w d x+\int_{\partial \Omega} b\left(x, u_{*}\right) w d \sigma=\int_{\Omega} h\left(x, u_{*}^{+}\right) w d x \quad \forall w \in X \tag{3.6}
\end{equation*}
$$

Put $w:=u_{*}^{-}$in (3.6) and exploit $\left(\mathrm{i}_{3}\right)$ of Lemma 2.1 to arrive at

$$
\frac{c_{3}}{p-1}\left\|\nabla u_{*}^{-}\right\|_{p}^{p}+\left\|u_{*}^{-}\right\|_{p}^{p} \leq 0
$$

Therefore, $u_{*} \geq 0$. Now, if $w:=\left(u_{*}-\bar{u}\right)^{+}$then (3.6), together with (3.1), $\left(\mathrm{f}_{1}\right)$, and (2.7), produce

$$
\begin{aligned}
\left\langle A\left(u_{*}\right),\right. & \left.\left(u_{*}-\bar{u}\right)^{+}\right\rangle+\int_{\Omega}(|\alpha|+1) u_{*}^{p-1}\left(u_{*}-\bar{u}\right)^{+} d x+\int_{\partial \Omega} \beta \bar{u}^{p-1}\left(u_{*}-\bar{u}\right)^{+} d \sigma \\
& =\int_{\Omega}\left(\eta \bar{u}^{q-1}-C_{\eta} \bar{u}^{r-1}+\bar{u}^{p-1}\right)\left(u_{*}-\bar{u}\right)^{+} d x \\
& \leq \int_{\Omega}\left[f(x, \bar{u})+\bar{u}^{p-1}\right]\left(u_{*}-\bar{u}\right)^{+} d x \\
& \leq\left\langle A(\bar{u}),\left(u_{*}-\bar{u}\right)^{+}\right\rangle+\int_{\Omega}(|\alpha|+1) \bar{u}^{p-1}\left(u_{*}-\bar{u}\right)^{+} d x+\int_{\partial \Omega} \beta \bar{u}^{p-1}\left(u_{*}-\bar{u}\right)^{+} d \sigma .
\end{aligned}
$$

This forces

$$
\left\langle A\left(u_{*}\right)-A(\bar{u}),\left(u_{*}-\bar{u}\right)^{+}\right\rangle+\int_{\Omega}(|\alpha|+1)\left(u_{*}^{p-1}-\bar{u}^{p-1}\right)\left(u_{*}-\bar{u}\right)^{+} d x \leq 0
$$

i.e., $u_{*} \leq \bar{u}$. Summing up, both $u_{*} \in[0, \bar{u}] \backslash\{0\}$ and, by (3.6) again,

$$
\begin{equation*}
\left\langle A\left(u_{*}\right), w\right\rangle+\int_{\Omega}|\alpha| u_{*}^{p-1} w d x+\int_{\partial \Omega} \beta u_{*}^{p-1} w d \sigma=\int_{\Omega}\left(\eta u_{*}^{q-1}-C_{\eta} u_{*}^{r-1}\right) w d x \forall w \in X \tag{3.7}
\end{equation*}
$$

Proposition 7 in [23] ensures that $u_{*} \in L^{\infty}(\Omega)$, whence $u_{*} \in C_{+} \backslash\{0\}$ thanks to Lieberman's regularity results [12]. Let $\alpha_{\theta} \in \mathbb{R}^{+}$satisfy

$$
\begin{equation*}
\eta t^{q-1}-C_{\eta} t^{r-1} \geq-\alpha_{\theta} t^{p-1}, \quad t \in[0, \theta] . \tag{3.8}
\end{equation*}
$$

Because of Remark 2.3, from (3.7)-(3.8) it follows

$$
\operatorname{div} a\left(\nabla u_{*}(x)\right) \leq\left(\|\alpha\|_{\infty}+\alpha_{\theta}\right) u_{*}(x)^{p-1} \text { a.e. in } \Omega .
$$

Thus, Pucci-Serrin's maximum principle [26, p. 120] yields $u_{*} \in \operatorname{int}\left(C_{+}\right)$.
Let us now come to uniqueness. Suppose $\hat{u} \in[0, \bar{u}] \cap \operatorname{int}\left(C_{+}\right)$is another solution to (3.4). Define, provided $u \in L^{1}(\Omega)$,

$$
J(u):= \begin{cases}\int_{\Omega} G\left(\nabla u^{\frac{1}{q}}\right) d x+\frac{1}{p}\left(\int_{\Omega}|\alpha| u^{\frac{p}{q}} d x+\int_{\partial \Omega} \beta u^{\frac{p}{q}} d \sigma\right) & \text { if } u \geq 0, u^{\frac{1}{q}} \in X, \\ +\infty & \text { otherwise }\end{cases}
$$

The reasoning made in [25, pp. 1219-1220] shows here that $u \mapsto \int_{\Omega} G\left(\nabla u^{\frac{1}{q}}\right) d x$ turns out convex. Since $p \geq q$ and $\beta \geq 0$, the same holds for $J$. Via Fatou's lemma we see that $J$ is lower semicontinuous. A simple computation chiefly based on [4, Theorem 2.4.54] gives

$$
\begin{aligned}
J^{\prime}\left(u_{*}^{q}\right)(w) & =\frac{1}{q} \int_{\Omega} \frac{-\operatorname{div} a\left(\nabla u_{*}\right)+|\alpha| u_{*}^{p-1}}{u_{*}^{q-1}} w d x \\
J^{\prime}\left(\hat{u}^{q}\right)(w) & =\frac{1}{q} \int_{\Omega} \frac{-\operatorname{div} a(\nabla \hat{u})+|\alpha| \hat{u}^{p-1}}{\hat{u}^{q-1}} w d x
\end{aligned}
$$

for all $w \in C^{1}(\bar{\Omega})$ (which is dense in $X$ ), while the monotonicity of $J^{\prime}$ entails

$$
\begin{aligned}
0 & \leq \int_{\Omega}\left(\frac{-\operatorname{div} a\left(\nabla u_{*}\right)+|\alpha| u_{*}^{p-1}}{u_{*}^{q-1}}-\frac{-\operatorname{div} a(\nabla \hat{u})+|\alpha| \hat{u}^{p-1}}{\hat{u}^{q-1}}\right)\left(u_{*}^{q}-\hat{u}^{q}\right) d x \\
& =\int_{\Omega}\left(\frac{\eta u_{*}^{q-1}-C_{\eta} u_{*}^{r-1}}{u_{*}^{q-1}}-\frac{\eta \hat{u}^{q-1}-C_{\eta} \hat{u}^{r-1}}{\hat{u}^{q-1}}\right)\left(u_{*}^{q}-\hat{u}^{q}\right) d x \\
& =-C_{\eta} \int_{\Omega}\left(u_{*}^{r-q}-\hat{u}^{r-q}\right)\left(u_{*}^{q}-\hat{u}^{q}\right) d x \leq 0
\end{aligned}
$$

as $q<r$. Consequently, $u_{*}=\hat{u}$. Working similarly produces a solution $v_{*}$ to (3.4) with the asserted properties.

Now, consider the sets

$$
\begin{aligned}
& \Sigma_{+}:=\{u \in X \backslash\{0\}: 0 \leq u \leq \bar{u}, u \text { solves }(1.1)\} \\
& \Sigma_{-}:=\{u \in X \backslash\{0\}: \underline{u} \leq u \leq 0, u \text { solves }(1.1)\} .
\end{aligned}
$$

Standard arguments show that:

- $\Sigma_{+} \subseteq \operatorname{int}\left(C_{+}\right)$while $\Sigma_{-} \subseteq-\operatorname{int}\left(C_{+}\right)$(cf. for instance the above proof);
- $\Sigma_{+}$is downward directed and $\Sigma_{-}$is upward directed (see, e.g., [24, Proposition 7]).

Lemma 3.2. Under the hypotheses of Lemma 3.1, one has

$$
u_{*} \leq u \quad \forall u \in \Sigma_{+}, \quad u \leq v_{*} \quad \forall u \in \Sigma_{-} .
$$

Proof. Pick any $u \in \Sigma_{+}$. Bearing in mind (3.2)-(3.3), define

$$
\begin{aligned}
& h_{+}(x, t):=\left\{\begin{array}{ll}
h\left(x, t^{+}\right) & \text {if } t \leq u(x) \\
h(x, u(x)) & \text { otherwise }
\end{array}, \quad H_{+}(x, t):=\int_{0}^{t} h_{+}(x, s) d s, \quad(x, t) \in \Omega \times \mathbb{R},\right. \\
& b_{+}(x, t):=\left\{\begin{array}{ll}
b\left(x, t^{+}\right) & \text {if } t \leq u(x) \\
b(x, u(x)) & \text { otherwise }
\end{array}, \quad B_{+}(x, t):=\int_{0}^{t} b_{+}(x, s) d s, \quad(x, t) \in \partial \Omega \times \mathbb{R} .\right.
\end{aligned}
$$

The associated functional

$$
\psi_{+}(w):=\int_{\Omega} G(\nabla w) d x+\frac{1}{p} \int_{\Omega}(|\alpha|+1)|w|^{p} d x+\int_{\partial \Omega} B_{+}(x, w) d \sigma-\int_{\Omega} H_{+}(x, w) d x
$$

is evidently $C^{1}$, weakly sequentially lower semi-continuous, and coercive. So, there exists $u_{0} \in X$ such that

$$
\psi_{+}\left(u_{0}\right)=\inf _{w \in X} \psi_{+}(w)
$$

From $q \leq p<r$ it follows, as before, $\psi_{+}\left(u_{0}\right)<0=\psi_{+}(0)$, namely $u_{0} \neq 0$. Moreover, $u_{0} \in[0, u]$, which entails

$$
\left\langle A\left(u_{0}\right), w\right\rangle+\int_{\Omega}(|\alpha|+1) u_{0}^{p-1} w d x+\int_{\partial \Omega} b\left(x, u_{0}\right) w d \sigma=\int_{\Omega} h\left(x, u_{0}\right) w d x, \quad w \in X .
$$

Through Lemma 3.1 we thus have $u_{0}=u_{*}$ and, a fortiori, $u_{*} \leq u$. The remaining proof is analogous.

Proposition 3.1. If $\left(\mathrm{a}_{1}\right)-\left(\mathrm{a}_{4}\right)$, (2.7), and $\left(\mathrm{f}_{1}\right)-\left(\mathrm{f}_{3}\right)$ hold then, there exists $u_{+} \in \Sigma_{+}$(resp., $v_{-} \in \Sigma_{-}$) such that $u_{+} \leq u$ for all $u \in \Sigma_{+}$(resp., $u \leq v_{-}$for all $u \in \Sigma_{-}$).

Proof. Both arguments are similar. Hence, we shall only present those involving $u_{+}$. Since $\Sigma_{+}$is downward directed, [10, Lemma 3.10] gives a sequence $\left\{u_{n}\right\} \subseteq \Sigma_{+}$, which fulfills

$$
u_{n+1} \leq u_{n} \quad \forall n \in \mathbb{N}, \quad \inf _{n \in \mathbb{N}} u_{n}=\inf \Sigma_{+}
$$

Consequently, $0 \leq u_{n} \leq\left\|u_{1}\right\|_{\infty}$ besides

$$
\begin{equation*}
\left\langle A\left(u_{n}\right), w\right\rangle+\int_{\Omega}|\alpha| u_{n}^{p-1} w d x+\int_{\partial \Omega} \beta u_{n}^{p-1} w d \sigma=\int_{\Omega} f\left(x, u_{n}\right) w d x, \quad w \in X \tag{3.9}
\end{equation*}
$$

for every $n \in \mathbb{N}$. Now, put $w:=u_{n}$ in (3.9) and exploit $\left(\mathrm{i}_{3}\right)$ of Lemma 2.1 to verify that $\left\{u_{n}\right\} \subseteq X$ turns out bounded. Let $u_{+} \in X_{+}$satisfy

$$
\begin{equation*}
u_{n} \rightharpoonup u_{+} \text {in } X, u_{n} \rightarrow u_{+} \text {in both } L^{p}(\Omega) \text { and } L^{p}(\partial \Omega), \tag{3.10}
\end{equation*}
$$

where a subsequence is considered if necessary. Combining (3.9) written for $w:=u_{n}-u_{+}$ with (3.10) entails

$$
\lim _{n \rightarrow+\infty}\left\langle A\left(u_{n}\right), u_{n}-u_{+}\right\rangle=0
$$

whence $u_{n} \rightarrow u_{+}$in $X$, because $A$ enjoys the (S) $)_{+}$property. Due to (3.9) again, this ensures that $u_{+}$solves (1.1). On the other hand, by Lemma 3.2, from $\left\{u_{n}\right\} \subseteq \Sigma_{+}$it follows $u_{*} \leq u_{n}$ for all $n \in \mathbb{N}$. Hence, $u_{*} \leq u_{+}$and, a fortiori, $u_{+} \in \Sigma_{+}$. Noting that $u_{+}=\inf _{n \in \mathbb{N}} u_{n}$ completes the proof.
Remark 3.3. On account of Proposition 3.1, every solution $u \in\left[v_{-}, u_{+}\right] \backslash\left\{v_{-}, 0, u_{+}\right\}$of (1.1) must be nodal.

For $u_{+}, v_{-}$as above and $\hat{\alpha}>\|\alpha\|_{\infty}$, define

$$
\begin{gather*}
\hat{f}(x, t):=\left\{\begin{array}{ll}
f\left(x, v_{-}(x)\right)+\hat{\alpha}\left|v_{-}(x)\right|^{p-2} v_{-}(x) & \text { when } t<v_{-}(x), \\
f(x, t)+\hat{\alpha}|t|^{p-2} t & \text { if } v_{-}(x) \leq t \leq u_{+}(x), \\
f\left(x, u_{+}(x)\right)+\hat{\alpha} u_{+}(x)^{p-1} & \text { when } t>u_{+}(x),
\end{array} \hat{F(x, t):=\int_{0}^{t} \hat{f}(x, s) d s,}\right. \tag{3.11}
\end{gather*}
$$

provided $(x, t) \in \Omega \times \mathbb{R}$, besides

$$
\begin{gather*}
\hat{b}(x, t):= \begin{cases}\beta(x)\left|v_{-}(x)\right|^{p-2} v_{-}(x) & \text { for } t<v_{-}(x), \\
\beta(x)|t|^{p-2} t & \text { if } v_{-}(x) \leq t \leq u_{+}(x), \\
\beta(x) u_{+}(x)^{p-1} & \text { for } t>u_{+}(x),\end{cases}  \tag{3.12}\\
\hat{B}(x, t):=\int_{0}^{t} \hat{b}(x, s) d s,
\end{gather*}
$$

where $(x, t) \in \partial \Omega \times \mathbb{R}$. A standard computation, which exploits ( $\mathrm{i}_{3}$ ) in Lemma 2.1, the choice of $\hat{\alpha}$, and (3.11)-(3.12), guarantees that the $C^{1}$-functionals

$$
\begin{aligned}
\hat{\varphi}(u) & :=\int_{\Omega} G(\nabla u) d x+\frac{1}{p} \int_{\Omega}(\alpha+\hat{\alpha})|u|^{p} d x+\int_{\partial \Omega} \hat{B}(x, u) d \sigma-\int_{\Omega} \hat{F}(x, u) d x, u \in X, \\
\hat{\varphi}_{ \pm}(u) & :=\int_{\Omega} G(\nabla u) d x+\frac{1}{p} \int_{\Omega}(\alpha+\hat{\alpha})|u|^{p} d x+\int_{\partial \Omega} \hat{B}(x, u) d \sigma-\int_{\Omega} \hat{F}\left(x, u^{ \pm}\right) d x, u \in X,
\end{aligned}
$$

are coercive; so, by Proposition 2.1, they comply with condition (PS). Moreover,
Lemma 3.3. Let $\left(\mathrm{a}_{1}\right)-\left(\mathrm{a}_{4}\right)$, (2.7), and $\left(\mathrm{f}_{1}\right)-\left(\mathrm{f}_{3}\right)$ be satisfied. Then:
$\left(\mathrm{j}_{1}\right) K(\hat{\varphi}) \subseteq\left[v_{-}, u_{+}\right] \cap C^{1}(\bar{\Omega})$.
$\left(\mathrm{j}_{2}\right) u_{+}$and $v_{-}$are local minimizers of $\hat{\varphi}$.
$\left(\mathrm{j}_{3}\right) K\left(\hat{\varphi}_{+}\right)=\left\{0, u_{+}\right\}$and $K\left(\hat{\varphi}_{-}\right)=\left\{0, v_{-}\right\}$.
Proof. Reasoning as before (cf. Lemma 3.1) we can easily check ( $\mathrm{j}_{1}$ ). A known argument (see., e.g., [17, Lemma 3.2] or [25, p. 1227, Claim 2]), chiefly based on Proposition 2.2, yields $\left(\mathrm{j}_{2}\right)$. Finally, concerning $\left(\mathrm{j}_{3}\right)$, let us simply note that the obvious inclusion $K\left(\hat{\varphi}_{+}\right) \subseteq\left[0, u_{+}\right]$forces $K\left(\hat{\varphi}_{+}\right)=\left\{0, u_{+}\right\}$by extremality of $u_{+}$; cf. Proposition 3.1. The same goes for $K\left(\hat{\varphi}_{-}\right)=\left\{0, v_{-}\right\}$.

Lemma 3.4. Under $\left(\mathrm{f}_{2}\right)-\left(\mathrm{f}_{4}\right)$ one has $C_{k}(\hat{\varphi}, 0)=0$ for all $k \in \mathbb{N}_{0}$.
Proof. It is rather delicate, but essentially analogous to the one made in [25, Proposition 4.1]. We shall present here a simpler trick. Fix $r>p$ and $\eta>0$. Assumptions $\left(\mathrm{f}_{2}\right)-\left(\mathrm{f}_{3}\right)$ give $C_{1}>0$ such that

$$
\hat{F}(x, t) \geq \eta|t|^{q}-C_{1}|t|^{r}, \quad(x, t) \in \Omega \times \mathbb{R} .
$$

Because of (2.5) this implies

$$
\hat{\varphi}(t u) \leq c_{9}\left(t^{q}\|\nabla u\|_{q}^{q}+t^{p}\|\nabla u\|_{p}^{p}\right)+C_{2} t^{p}\|u\|_{p}^{p}+C_{1} t^{r}\|u\|_{r}^{r}-\eta t^{q}\|u\|_{q}^{q}
$$

for every $t>0, u \in X$. Since $\eta$ was arbitrary while $q \leq p<r$, if $u \neq 0$ then there exists $\left.t^{*} \in\right] 0,1[($ which may depend on $u)$ fulfilling $\hat{\varphi}(t u)<0$ whatever $t \in] 0, t^{*}[$. Define $t_{1}:=\sup \{t \in[0,1]: \hat{\varphi}(t u)<0\}$ as well as

$$
t_{2}:= \begin{cases}\inf \{t \in[0,1]: \hat{\varphi}(t u) \geq 0\} & \text { when }\{t \in[0,1]: \hat{\varphi}(t u) \geq 0\} \neq \emptyset \\ 1 & \text { otherwise }\end{cases}
$$

We will show that $t_{1} \leq t_{2}$. By contradiction, suppose $t_{2}<t_{1}$. Let $\left.\left.t_{0} \in\right] 0,1\right]$ satisfy $\hat{\varphi}(v)=0$, where $v:=t_{0} u$. Simple calculations based on $\left(\mathrm{f}_{2}\right)-\left(\mathrm{f}_{4}\right)$ and (3.1) yield

$$
\hat{q} F(x, t)-f(x, t) t \geq\left[\left(1-\frac{\hat{q}}{p}\right)\|\alpha\|_{\infty}+k\right]|t|^{p}-C_{3}|t|^{r}, \quad(x, t) \in \Omega \times[-\theta, \theta],
$$

for some $k>0, C_{3}>0$. Thanks to ( $\mathrm{a}_{4}$ ), this inequality entails

$$
\begin{aligned}
t_{0} \frac{d}{d t} \hat{\varphi}(t u) L_{t=t_{0}} & =\left\langle\hat{\varphi}^{\prime}(v), v\right\rangle=\left\langle\hat{\varphi}^{\prime}(v), v\right\rangle-\hat{q} \hat{\varphi}(v) \\
& \geq c_{6}\|\nabla v\|_{p}^{p}+\left(1-\frac{\hat{q}}{p}\right) \int_{\Omega} \alpha|v|^{p} d x+\int_{\Omega}(\hat{q} F(x, v)-f(x, v) v) d x \\
& \geq c_{6}\|\nabla v\|_{p}^{p}+k\|v\|_{p}^{p}-C_{3}\|v\|_{r}^{r}>0
\end{aligned}
$$

provided $0<\|u\| \leq \rho$, with $\rho$ small enough. Thus,

$$
\begin{equation*}
\left.\frac{d}{d t} \hat{\varphi}(t u)\left\lfloor_{t=t_{0}}>0 \text { whenever } t_{0} \in\right] 0,1\right], \hat{\varphi}\left(t_{0} u\right)=0 \tag{3.13}
\end{equation*}
$$

Since $\hat{\varphi}\left(t_{2} u\right)=0$, from (3.13) it follows

$$
\begin{equation*}
\left.\hat{\varphi}(t u)>0 \quad \forall t \in] t_{2}, t_{2}+\delta_{1}\right], \tag{3.14}
\end{equation*}
$$

where $0<\delta_{1}<t_{1}-t_{2}$. Letting

$$
\hat{t}:= \begin{cases}\min \left\{t \in\left[t_{2}+\delta_{1}, t_{1}\right]: \hat{\varphi}(t u)=0\right\} & \text { if }\left\{t \in\left[t_{2}+\delta_{1}, t_{1}\right]: \hat{\varphi}(t u)=0\right\} \neq \emptyset \\ 1 & \text { otherwise },\end{cases}
$$

(3.14) forces $\hat{t}>t_{2}+\delta_{1}$. So, via (3.13) when $\hat{\varphi}(\hat{t} u)=0$, we can find a $\left.\delta_{2} \in\right] 0, \hat{t}-t_{2}-\delta_{1}[$ such that

$$
\begin{equation*}
\hat{\varphi}(t u)<0, \quad t \in\left[\hat{t}-\delta_{2}, \hat{t}[.\right. \tag{3.15}
\end{equation*}
$$

Now, by (3.14), (3.15), and Bolzano's theorem, $\hat{\varphi}\left(t^{*} u\right)=0$ for some $\left.t^{*} \in\right] t_{2}+\delta_{1}, \hat{t}-\delta_{2}[$, which is impossible due to the choice of $\hat{t}$. Therefore, $t_{1} \leq t_{2}$, as desired.
One actually has $t_{1}=t_{2}$, because assuming $t_{1}<t_{2}$ leads to $\hat{\varphi}(t u)=0$ in $] t_{1}, t_{2}[$, against (3.13). Put $t(u):=t_{1}=t_{2}$. Evidently,

$$
\hat{\varphi}(t u)<0 \quad \forall t \in] 0, t(u)[, \quad \hat{\varphi}(t(u) u)=0, \quad \hat{\varphi}(t u)>0 \quad \forall t \in] t(u), 1],
$$

whence the map $r(u):=t(u) u, u \in \bar{B}_{\rho} \backslash\{0\}$, turns out continuous,

$$
r\left(\bar{B}_{\rho} \backslash\{0\}\right) \subseteq\left(\hat{\varphi}^{0} \cap \bar{B}_{\rho}\right) \backslash\{0\}, \quad r L_{\left(\hat{\varphi}^{0} \cap \bar{B}_{\rho}\right) \backslash\{0\}}=i d\left\lfloor_{\left(\hat{\varphi}^{0} \cap \bar{B}_{\rho}\right) \backslash\{0\}} .\right.
$$

This shows that $\left(\hat{\varphi}^{0} \cap \bar{B}_{\rho}\right) \backslash\{0\}$ is a retract of $\bar{B}_{\rho} \backslash\{0\}$. Consequently, $\left(\hat{\varphi}^{0} \cap \bar{B}_{\rho}\right) \backslash\{0\}$ turns out contractible in itself, since $\bar{B}_{\rho} \backslash\{0\}$ enjoys the same property. Now, Propositions 4.9-4.10 of [7] give

$$
H_{k}\left(\hat{\varphi}^{0} \cap \bar{B}_{\rho},\left(\hat{\varphi}^{0} \cap \bar{B}_{\rho}\right) \backslash\{0\}\right)=0, \quad k \in \mathbb{N}_{0}
$$

i.e., the conclusion.

We are now ready to establish our first existence result.
Theorem 3.1. If $\left(\mathrm{a}_{1}\right)-\left(\mathrm{a}_{4}\right),(2.7)$, and $\left(\mathrm{f}_{1}\right)-\left(\mathrm{f}_{4}\right)$ hold true then (1.1) admits a nodal solution $\hat{u} \in C^{1}(\bar{\Omega}) \cap\left[v_{-}, u_{+}\right]$.

Proof. Since $\hat{\varphi}$ turns out coercive, via [19, Proposition 6.64] one has $C_{k}(\hat{\varphi}, \infty)=\delta_{k, 0} \mathbb{Z}$. Combining ( $\mathrm{j}_{2}$ ) of Lemma 3.3 with [19, Example 6.45] next entails

$$
\begin{equation*}
C_{k}\left(\hat{\varphi}, u_{+}\right)=C_{k}\left(\hat{\varphi}, v_{-}\right)=\delta_{k, 0} \mathbb{Z} \tag{3.16}
\end{equation*}
$$

Suppose $K(\hat{\varphi})=\left\{0, u_{+}, v_{-}\right\}$, recall Lemma 3.4, and write the Morse relation (2.1) for $t:=-1$, to arrive at $2(-1)^{0}=(-1)^{0}$, which is evidently impossible. Thus, there exists a point $\hat{u} \in K(\hat{\varphi}) \backslash\left\{0, u_{+}, v_{-}\right\}$. The conclusion easily stems from ( $\mathrm{j}_{1}$ ) of Lemma 3.3 besides (3.11)-(3.12).

### 3.2 The $(q-1)$-linear case

For $q, c_{7}$ given by $\left(\mathrm{a}_{4}\right), \alpha_{0}:=|\alpha| / c_{7}$, and $\beta_{0}:=\beta / c_{7}$, assume that:
( $\mathrm{f}_{5}$ ) Uniformly in $x \in \Omega$, one has

$$
c_{7} \hat{\lambda}_{2}\left(q, \alpha_{0}, \beta_{0}\right)<c_{10}<\liminf _{t \rightarrow 0} \frac{f(x, t)}{|t|^{q-2} t} \leq \limsup _{t \rightarrow 0} \frac{f(x, t)}{|t|^{q-2} t} \leq c_{11} .
$$

A careful inspection of proofs reveals that all the auxiliary results above, except Lemma 3.4, remain valid whenever $\left(f_{5}\right)$ replaces $\left(f_{3}\right)$. So, although critical groups cannot be employed, the same conclusion is achieved via ( $p_{3}$ ) in Section 2.

Theorem 3.2. Let $\left(\mathrm{a}_{1}\right)-\left(\mathrm{a}_{4}\right)$, (2.7), $\left(\mathrm{f}_{1}\right)-\left(\mathrm{f}_{2}\right)$, and $\left(\mathrm{f}_{5}\right)$ be satisfied. Then (1.1) possesses a nodal solution $\hat{u} \in C^{1}(\bar{\Omega})$.

Proof. Recalling $\left(\mathrm{j}_{1}\right)-\left(\mathrm{j}_{2}\right)$ of Lemma 3.3, we may suppose $K(\hat{\varphi})$ finite, the local minimizer $v_{-}, u_{+}$proper, besides $\hat{\varphi}\left(v_{-}\right) \leq \hat{\varphi}\left(u_{+}\right)$(the other case is analogous). If $0<\rho<\left\|u_{+}-v_{-}\right\|$ complies with

$$
\begin{equation*}
\hat{\varphi}\left(u_{+}\right)<C_{\rho}:=\inf _{u \in \partial B_{\rho}\left(u_{+}\right)} \hat{\varphi}(u) \tag{3.17}
\end{equation*}
$$

then the Mountain Pass theorem produces a point $u_{1} \in K(\hat{\varphi})$ such that

$$
\begin{equation*}
C_{\rho} \leq \hat{\varphi}\left(u_{1}\right)=\inf _{\gamma \in \Gamma} \max _{t \in[0,1]} \hat{\varphi}(\gamma(t)), \tag{3.18}
\end{equation*}
$$

where

$$
\Gamma:=\left\{\gamma \in C^{0}([0,1], X): \gamma(0)=v_{-}, \gamma(1)=u_{+}\right\} .
$$

By ( $\mathrm{j}_{1}$ ) in Lemma 3.3, $u_{1}$ belongs to $C^{1}(\bar{\Omega})$ and solves (1.1). Through (3.17)-(3.18) we next get $u_{1} \neq v_{-}, u_{+}$. Thus, on account of Proposition 3.1, it remains to check whether $u_{1} \neq 0$. This will follow from the inequality $\hat{\varphi}\left(u_{1}\right)<0$, which evidently holds once there exists a path $\tilde{\gamma} \in \Gamma$ such that

$$
\begin{equation*}
\hat{\varphi}(\tilde{\gamma}(t))<0 \quad \forall t \in[0,1] . \tag{3.19}
\end{equation*}
$$

Fix $\varepsilon>0$. Using ( $\mathrm{a}_{4}$ ) yields

$$
\begin{equation*}
G(\xi) \leq \frac{c_{7}+\varepsilon}{q}|\xi|^{q}, \quad|\xi| \leq \delta \tag{3.20}
\end{equation*}
$$

while ( $\mathrm{f}_{5}$ ) entails

$$
\begin{equation*}
F(x, t) \geq \frac{c_{10}}{q}|t|^{q}, \quad(x, t) \in \Omega \times[-\delta, \delta], \tag{3.21}
\end{equation*}
$$

provided $\delta>0$ is small enough. Now, denote by $\mathcal{E}_{q}^{0}$ the functional (2.8) with $\alpha_{0}$ and $\beta_{0}$ in place of $\alpha$ and $\beta$, respectively. Thanks to ( $\mathrm{p}_{3}$ ) besides Lemma 2.3, given $\eta>0$ we can find a path $\hat{\gamma}_{\eta} \in \hat{\Gamma}_{C}$ such that

$$
\begin{equation*}
\max _{t \in[-1,1]} c_{7} \mathcal{E}_{q}^{0}\left(\hat{\gamma}_{\eta}(t)\right)<c_{7} \hat{\lambda}_{2}\left(q, \alpha_{0}, \beta_{0}\right)+\eta . \tag{3.22}
\end{equation*}
$$

Since $\hat{\gamma}_{\eta}([-1,1])$ is compact in $C^{1}(\bar{\Omega})$ while $-v_{-}, u_{+} \in \operatorname{int}\left(C_{+}\right)$, there exists $\tau>0$ fulfilling

$$
v_{-}(x) \leq \tau \hat{\gamma}_{\eta}(t)(x) \leq u_{+}(x), \quad\left|\tau \hat{\gamma}_{\eta}(t)(x)\right| \leq \delta \leq 1, \quad\left|\tau \nabla \hat{\gamma}_{\eta}(t)(x)\right| \leq \delta
$$

for all $(x, t) \in \Omega \times[-1,1]$. On account of $q \leq p$, the inequalities above, (3.20)-(3.22), and $\left\|\hat{\gamma}_{\eta}(t)\right\|_{q} \equiv 1$, one arrives at

$$
\hat{\varphi}\left(\tau \hat{\gamma}_{\eta}(t)\right) \leq \frac{\tau^{q}}{q}\left[c_{7} \mathcal{E}_{q}^{0}\left(\hat{\gamma}_{\eta}(t)\right)+\varepsilon\left\|\nabla \hat{\gamma}_{\eta}(t)\right\|_{q}^{q}-c_{10}\right]<\frac{\tau^{q}}{q}\left[c_{7} \hat{\lambda}_{2}\left(q, \alpha_{0}, \beta_{0}\right)+\eta+C_{1} \varepsilon-c_{10}\right] .
$$

Therefore,

$$
\begin{equation*}
\hat{\varphi}\left(\tau \hat{\gamma}_{n}(t)\right)<0 \quad \forall t \in[-1,1] \tag{3.23}
\end{equation*}
$$

as soon as $\varepsilon$ and $\eta$ are taken so small that

$$
\eta+C_{1} \varepsilon<c_{10}-c_{7} \hat{\lambda}_{2}\left(q, \alpha_{0}, \beta_{0}\right)
$$

see $\left(f_{5}\right)$. Next, if $\hat{a}:=\hat{\varphi}_{+}\left(u_{+}\right)$then $\hat{a}<0$, because

$$
\hat{\varphi}_{+}\left(u_{+}\right)=\inf _{u \in X} \hat{\varphi}_{+}(u)<0=\hat{\varphi}_{+}(0),
$$

no critical value of $\hat{\varphi}_{+}$lies in $(\hat{a}, 0)$, and $K_{\hat{a}}\left(\hat{\varphi}_{+}\right)=\left\{u_{+}\right\}$; see Lemma 3.3. Thus, the second deformation lemma [4, Theorem 5.1.33] gives a continuous map $h:[0,1] \times\left(\hat{\varphi}_{+}^{0} \backslash\{0\}\right) \rightarrow \hat{\varphi}_{+}^{0}$ satisfying

$$
\begin{equation*}
h(0, u)=u, \quad h(1, u)=u_{+}, \quad \hat{\varphi}_{+}(h(t, u)) \leq \hat{\varphi}_{+}(u) \tag{3.24}
\end{equation*}
$$

for all $(t, u) \in[0,1] \times\left(\hat{\varphi}_{+}^{0} \backslash\{0\}\right)$. By (3.23) one has

$$
\hat{\varphi}_{+}\left(\tau \hat{u}_{1}\left(q, \alpha_{0}, \beta_{0}\right)\right)=\hat{\varphi}\left(\tau \hat{\gamma}_{\eta}(1)\right)<0 .
$$

Hence, it makes sense to define

$$
\gamma_{+}(t):=h\left(t, \tau \hat{u}_{1}\left(q, \alpha_{0}, \beta_{0}\right)\right)^{+}, \quad t \in[0,1] .
$$

The path $\gamma_{+}:[0,1] \rightarrow X$ connects $\tau \hat{u}_{1}\left(q, \alpha_{0}, \beta_{0}\right)$ with $u_{+}$. Moreover, due to (3.23)-(3.24),

$$
\begin{equation*}
\hat{\varphi}\left(\gamma_{+}(t)\right)=\hat{\varphi}_{+}\left(\gamma_{+}(t)\right) \leq \hat{\varphi}_{+}\left(\tau \hat{u}_{1}\left(q, \alpha_{0}, \beta_{0}\right)\right)=\hat{\varphi}\left(\tau \hat{\gamma}_{\eta}(1)\right)<0 \quad \forall t \in[0,1] . \tag{3.25}
\end{equation*}
$$

A similar reasoning, where $\hat{\varphi}_{-}$takes the place of $\hat{\varphi}_{+}$, produces a continuous function $\gamma_{-}:[0,1] \rightarrow X$ such that $\gamma_{-}(0)=v_{-}, \gamma_{-}(1)=-\tau \hat{u}_{1}\left(q, \alpha_{0}, \beta_{0}\right)$, as well as

$$
\begin{equation*}
\hat{\varphi}\left(\gamma_{-}(t)\right)<0 \text { in }[0,1] . \tag{3.26}
\end{equation*}
$$

Concatenating $\gamma_{-}, \tau \hat{\gamma}_{\eta}$, and $\gamma_{+}$one obtains a path $\tilde{\gamma} \in \Gamma$ which, in view of (3.25)-(3.26), besides (3.23), fulfills (3.19).

### 3.3 The case of ( $p, 2$ )-Laplacian

Let $p>2, a_{0}(t):=t^{p-2}+1$, namely $q:=2, f(x, \cdot) \in C^{1}([-\theta, \theta])$ for every $x \in \Omega$, and let $\hat{\lambda}_{n}, E\left(\hat{\lambda}_{n}\right), \bar{H}_{n}, \hat{H}_{n}$ be like at the end of Section 2. The following assumptions will be posited.
( $\mathrm{f}_{6}$ ) There exists $a_{\theta} \in L^{\infty}(\Omega)$ such that $\left|f_{t}^{\prime}(x, t)\right| \leq a_{\theta}(x)$ in $\Omega \times[-\theta, \theta]$.
$\left(\mathrm{f}_{7}\right)$ With appropriate $m \geq 2, \delta_{0}>0$ small, $b \in L^{\infty}(\Omega) \backslash\left\{\hat{\lambda}_{m+1}\right\}$ one has $f(x, t) t \geq \hat{\lambda}_{m} t^{2}$ for all $(x, t) \in \Omega \times\left[-\delta_{0}, \delta_{0}\right]$ and

$$
f_{t}^{\prime}(x, 0)=\lim _{t \rightarrow 0} \frac{f(x, t)}{t} \leq b(x) \leq \hat{\lambda}_{m+1} \text { uniformly in } x \in \Omega
$$

Remark 3.4. As before, except Lemma 3.4, the auxiliary results above remain valid if $\left(\mathrm{f}_{6}\right)-\left(\mathrm{f}_{7}\right)$ replace $\left(\mathrm{f}_{3}\right)$.
Now, recall (3.11)-(3.12) and define

$$
\hat{\psi}(u):=\frac{1}{2}\left[\|\nabla u\|_{2}^{2}+\int_{\Omega}(\alpha+\hat{\alpha}) u^{2} d x\right]+\int_{\partial \Omega} \hat{B}(x, u) d \sigma-\int_{\Omega} \hat{F}(x, u) d x, u \in H^{1}(\Omega) .
$$

Evidently, $\hat{\psi}$ is $C^{2}$ in a neighborhood of the origin, besides $C^{1}$ on the whole $H^{1}(\Omega)$.
Lemma 3.5. Hypotheses (2.7) and $\left(\mathrm{f}_{6}\right)-\left(\mathrm{f}_{7}\right)$ entail $C_{k}(\hat{\psi}, 0)=\delta_{k, d_{m}} \mathbb{Z}$ for every $k \in \mathbb{N}_{0}$, where $d_{m}:=\operatorname{dim}\left(\bar{H}_{m}\right) \geq 2$.

Proof. Since $\bar{H}_{m}$ is finite dimensional, we can find $\rho_{1}>0$ such that

$$
u \in \bar{H}_{m} \cap \bar{B}_{\rho_{1}} \Longrightarrow|u(x)| \leq \delta_{0} \forall x \in \Omega
$$

Via $\left(\mathrm{f}_{7}\right)$ and $\left(\mathrm{p}_{4}\right)$, this easily leads to

$$
\begin{equation*}
\hat{\psi}(u) \leq \frac{1}{2}\left[\mathcal{E}_{2}(u)-\hat{\lambda}_{m}\|u\|_{2}^{2}\right] \leq 0, \quad u \in \bar{H}_{m} \cap \bar{B}_{\rho_{1}} \tag{3.27}
\end{equation*}
$$

Next, given $\varepsilon>0, r>2$, assumption $\left(\mathrm{f}_{7}\right)$ yields

$$
\hat{F}(x, t) \leq \frac{1}{2}(b(x)+\varepsilon) t^{2}+C_{1}|t|^{r} \text { in } \Omega \times \mathbb{R}
$$

whence, by [1, Lemma 2.2],

$$
\hat{\psi}(u) \geq \frac{1}{2}\left[\mathcal{E}_{2}(u)-\int_{\Omega} b(x) u^{2} d x-\varepsilon\|u\|^{2}\right]-C_{1}\|u\|^{r} \geq \frac{\hat{c}-\varepsilon}{2}\|u\|^{2}-C_{1}\|u\|^{r} \forall u \in \hat{H}_{m+1} .
$$

Here, $\|\cdot\|$ denotes the usual norm of $H^{1}(\Omega)$. Choosing $\varepsilon<\hat{c}$ we thus achieve

$$
\begin{equation*}
\hat{\psi}(u)>0, u \in \hat{H}_{m+1} \cap \bar{B}_{\rho_{2}} \backslash\{0\} \tag{3.28}
\end{equation*}
$$

provided $\rho_{2}>0$ is small enough. Inequalities (3.27)-(3.28) ensure that $\hat{\psi}$ has a local linking at zero with respect to the sum decomposition $H^{1}(\Omega)=\bar{H}_{m} \oplus V$, where $V$ indicates the closure of $\hat{H}_{m+1}$ in $H^{1}(\Omega)$. Since $\psi$ is coercive, it satisfies condition (PS); see Proposition 2.1. So, the conclusion follows from [27, Proposition 2.3].

Theorem 3.3. Let (2.7), $\left(\mathrm{f}_{1}\right)$, $\left(\mathrm{f}_{6}\right)$, and $\left(\mathrm{f}_{7}\right)$ be satisfied. Then (1.1), where $p>2$ while $a_{0}(t):=t^{p-2}+1$, admits a nodal solution $\hat{u} \in C^{1}(\bar{\Omega})$.

Proof. Set $\psi:=\hat{\psi}\left\lfloor_{X}\right.$. One evidently has $C_{k}(\psi, 0)=C_{k}(\hat{\psi}, 0)$, because $X \hookrightarrow H^{1}(\Omega)$ densely. Consequently, thanks to Lemma 3.5,

$$
\begin{equation*}
C_{k}(\psi, 0)=\delta_{k, d_{m}} \mathbb{Z} \tag{3.29}
\end{equation*}
$$

Observe next that

$$
|\hat{\varphi}(u)-\psi(u)| \leq C_{1}\|u\|^{p}, \quad\left|\left\langle\hat{\varphi}^{\prime}(u)-\psi^{\prime}(u), v\right\rangle\right| \leq C_{2}\|u\|^{p-1}\|v\| \quad \forall u, v \in X
$$

as a simple computation shows. Therefore, the $C^{1}$ - continuity of critical groups $[6$, Theorem 5.126] and (3.29) produce

$$
\begin{equation*}
C_{k}(\hat{\varphi}, 0)=\delta_{k, d_{m}} \mathbb{Z}, \quad k \in \mathbb{N}_{0} \tag{3.30}
\end{equation*}
$$

On the other hand, $\hat{\varphi}$ is coercive, whence $\inf _{u \in X} \hat{\varphi}(u)>-\infty$, and fulfills (PS). By [19, Proposition 6.64] we thus get

$$
\begin{equation*}
C_{k}(\hat{\varphi}, \infty)=\delta_{k, 0} \mathbb{Z}, \quad k \in \mathbb{N}_{0} . \tag{3.31}
\end{equation*}
$$

Combining (3.30)-(3.31) with [19, Corollary 6.92] one arrives at

$$
\begin{equation*}
C_{d_{m}-1}(\hat{\varphi}, \hat{u}) \neq 0 \text { or } C_{d_{m}+1}(\hat{\varphi}, \hat{u}) \neq 0, \text { where } d_{m} \geq 2 \tag{3.32}
\end{equation*}
$$

for some $\hat{u} \in K(\hat{\varphi}) \backslash\{0\}$. Now, the conclusion easily stems from (3.16), (3.32), besides ( $\mathrm{j}_{1}$ ) in Lemma 3.3; see also Remark 3.3.

## 4 Nodal solutions: multiplicity

Under a symmetry condition on $f(x, \cdot)$, problem (1.1) possesses infinitely many signchanging solutions.

Theorem 4.1. Suppose $\left(\mathrm{a}_{1}\right)-\left(\mathrm{a}_{4}\right),(2.7)$, and $\left(\mathrm{f}_{1}\right)-\left(\mathrm{f}_{3}\right)$ hold. If, moreover,
$\left(\mathrm{f}_{8}\right)$ the function $t \mapsto f(x, t)$ is odd in $[-\theta, \theta]$ for every $x \in \Omega$
then there exists a sequence $\left\{u_{n}\right\} \subseteq C^{1}(\bar{\Omega})$ of distinct nodal solutions to (1.1) satisfying $u_{n} \rightarrow 0$ in $C^{1}(\bar{\Omega})$.

Proof. The proof is patterned after that of [14, Theorem 4.3]; so, we only sketch it. Via $\left(\mathrm{a}_{4}\right)$ one has

$$
\begin{equation*}
G(\xi) \leq \frac{C_{1}}{q}|\xi|^{q}, \quad|\xi| \leq \delta, \tag{4.1}
\end{equation*}
$$

while, given $\eta>0$, assumption ( $\mathrm{f}_{3}$ ) entails

$$
\begin{equation*}
F(x, t) \geq \frac{\eta}{q}|t|^{q}, \quad(x, t) \in \Omega \times[-\delta, \delta], \tag{4.2}
\end{equation*}
$$

with $\delta>0$ small enough. Let $V \subseteq X$ be any finite dimensional subspace and let $\rho>0$ fulfill

$$
\begin{equation*}
u \in V \cap \bar{B}_{\rho} \quad \Longrightarrow \quad|u(x)| \leq \delta \forall x \in \Omega \tag{4.3}
\end{equation*}
$$

Gathering (4.1)-(4.3) together leads to

$$
\hat{\varphi}(u) \leq \frac{C_{1}}{q}\|\nabla u\|_{q}^{q}+\frac{1}{p}\left(\int_{\Omega}|\alpha||u|^{p} d x+\int_{\partial \Omega} \beta|u|^{p} d \sigma\right)-\frac{\eta}{q}\|u\|_{q}^{q} \leq\left(C_{2}-C_{3} \eta\right)\|u\|_{q}^{q}<0
$$

provided $u \in\left(V \cap \bar{B}_{\rho}\right) \backslash\{0\}, \eta>C_{2} / C_{3}$. Here, the equivalence between all norms on $V$ was also exploited. Hence, Theorem 1 of [11] furnishes a sequence

$$
\begin{equation*}
\left\{u_{n}\right\} \subseteq K(\hat{\varphi}) \cap\{u \in X: \hat{\varphi}(u)<0\} \tag{4.4}
\end{equation*}
$$

that converges to zero in $X$. Through standard arguments from the nonlinear regularity theory we actually have $\left\{u_{n}\right\} \subseteq C^{1}(\bar{\Omega})$ as well as $\left\|u_{n}\right\|_{C^{1}(\bar{\Omega})} \rightarrow 0$. Now, assertion ( $\mathrm{j}_{1}$ ) of Lemma 3.3, besides (4.4), easily yield the conclusion.

When the left-hand side is the ( $p, 2$ )-Laplacian, one can do without symmetry. However, a further condition on $f$ will be imposed.
( $\mathrm{f}_{9}$ ) There exists $\mu_{\theta}>0$ such that $t \mapsto f(x, t)+\mu_{\theta}|t|^{p-2} t$ is non-decreasing on $[-\theta, \theta]$ for all $x \in \Omega$.

Theorem 4.2. Let (2.7), $\left(\mathrm{f}_{1}\right)-\left(\mathrm{f}_{3}\right)$, and $\left(\mathrm{f}_{9}\right)$ be satisfied. If $p>2, a_{0}(t):=t^{p-2}+1$, while $\beta>0$ on $\partial \Omega$ then (1.1) possesses two nodal solutions $\hat{u}, \tilde{u} \in C^{1}(\bar{\Omega})$.

Proof. A first sign-changing function $\hat{u} \in C^{1}(\bar{\Omega}) \cap\left[v_{-}, u_{+}\right]$that solves (1.1) directly comes from Theorem 3.1. Since $p>2$ and $a(\xi)=\left(|\xi|^{p-2}+1\right) \xi$, an easy computation shows that

$$
(\nabla a(\xi) y) \cdot y \geq|y|^{2} \quad \forall \xi, y \in \mathbb{R}^{N} .
$$

Hence, the tangency principle [26, Theorem 2.5.2] gives

$$
\begin{equation*}
v_{-}<\hat{u}<u_{+} \text {in } \Omega . \tag{4.5}
\end{equation*}
$$

Now, pick any $\mu>\mu_{\theta}$ and define

$$
h_{1}:=f(\cdot, \hat{u})+\mu_{\theta}|\hat{u}|^{p-2} \hat{u}+\left(\mu-\mu_{\theta}\right)|\hat{u}|^{p-2} \hat{u}, \quad h_{2}:=f\left(\cdot, u_{+}\right)+\mu_{\theta} u_{+}^{p-1}+\left(\mu-\mu_{\theta}\right) u_{+}^{p-1} .
$$

Via ( $\mathrm{f}_{1}$ ) we obtain $h_{1}, h_{2} \in L^{\infty}(\Omega)$ while ( $\mathrm{f}_{9}$ ) entails $h_{1} \leq h_{2}$. Thanks to (4.5), for every compact set $K \subseteq \Omega$ one has $\operatorname{ess}_{\inf }^{x \in K}$ ( $\left.h_{2}(x)-h_{1}(x)\right)>0$. The condition on $\beta$ forces

$$
\frac{\partial u_{+}}{\partial n_{a}} L_{\partial \Omega}=-\beta u_{+}^{p-1}<0 .
$$

Consequently, by [2, Proposition 3], $u_{+}-\hat{u} \in \operatorname{int}\left(C_{+}\right)$. A quite similar reasoning produces $\hat{u}-v_{-} \in \operatorname{int}\left(C_{+}\right)$, whence, a fortiori, $\hat{u} \in \operatorname{int}_{C^{1}(\bar{\Omega})}\left(\left[v_{-}, u_{+}\right]\right)$. At this point, we adapt the flow invariance arguments made in $[8]$ to get a nodal solution $\tilde{u} \in C^{1}(\bar{\Omega}) \backslash \operatorname{int}_{C^{1}(\bar{\Omega})}\left(\left[v_{-}, u_{+}\right]\right)$ of (1.1). It is evident that $\tilde{u} \neq \hat{u}$.

A better situation occurs in the semi-linear case $p:=2$ and $a_{0}(t):=1$, because the regularity theory of [28] allows to weaken (2.7) as follows.

$$
\begin{equation*}
a \in L^{s}(\Omega) \text { for some } s>N, a^{+} \in L^{\infty}(\Omega), \beta \in W^{2, \infty}(\Omega), \text { and } \beta \geq 0 . \tag{4.6}
\end{equation*}
$$

Write $\hat{m}:=\max \left\{n_{0}, 2\right\}$, where $n_{0}:=\inf \left\{n \in \mathbb{N}: \hat{\lambda}_{n}>0\right\}$.
Theorem 4.3. If $p:=2, a_{0}(t):=1$, and (4.6), $\left(\mathrm{f}_{1}\right),\left(\mathrm{f}_{6}\right),\left(\mathrm{f}_{7}\right)$ with $m \geq \hat{m}$ hold true then (1.1) admits three nodal solutions $\hat{u}, \bar{u}, \tilde{u} \in C^{1}(\bar{\Omega})$.

Proof. The same technique exploited to prove both [16, Theorem 3.2] and [1, Theorem 3.2 ] yields a point $\hat{u} \in K(\hat{\varphi})$ of mountain pass type. So,

$$
\begin{equation*}
C_{k}(\hat{\varphi}, \hat{u})=\delta_{k, 1} \mathbb{Z} \quad \forall k \in \mathbb{N}_{0} ; \tag{4.7}
\end{equation*}
$$

vide [16, Lemma 3.1]. Recalling Remarks 3.3-3.4, conclusions ( $\mathrm{j}_{1}$ )- ( $\mathrm{j}_{2}$ ) in Lemma 3.3 ensure that $\hat{u} \in C^{1}(\bar{\Omega})$ is a nodal solution to (1.1). Via ( $\mathrm{f}_{6}$ ) we can find a $\mu_{\theta}>0$ such that $t \mapsto f(x, t)+\mu_{\theta} t$ turns out non-decreasing on $[-\theta, \theta]$ for any $x \in \Omega$. Since $\hat{u} \leq u_{+}$, one has

$$
\begin{array}{r}
-\Delta \hat{u}(x)+\left[a(x)+\mu_{\theta}\right] \hat{u}(x)=f(x, \hat{u}(x))+\mu_{\theta} \hat{u}(x) \\
\leq f\left(x, u_{+}(x)\right)+\mu_{\theta} u_{+}(x)=-\Delta u_{+}(x)+\left[a(x)+\mu_{\theta}\right] u_{+}(x),
\end{array}
$$

which entails

$$
\Delta\left(u_{+}-\hat{u}\right)(x) \leq\left(\left\|a^{+}\right\|_{\infty}+\mu_{\theta}\right)\left[u_{+}(x)-\hat{u}(x)\right], \quad x \in \Omega .
$$

From the strong maximum principle [26, p. 34] it follows $u_{+}-\hat{u}>0$ in $\Omega$. Suppose $\left(u_{+}-\hat{u}\right)\left(x_{0}\right)=0$ for some $x_{0} \in \partial \Omega$. The boundary point lemma [26, p. 120] leads to $\frac{\partial\left(u_{+}-\hat{u}\right)}{\partial n}\left(x_{0}\right)<0$, whence

$$
-\beta\left(x_{0}\right) u_{+}\left(x_{0}\right)=\frac{\partial u_{+}}{\partial n}\left(x_{0}\right)<\frac{\partial \hat{u}}{\partial n}\left(x_{0}\right)=-\beta\left(x_{0}\right) \hat{u}\left(x_{0}\right) .
$$

However, this is impossible, because $\beta \geq 0$. Thus, $u_{+}-\hat{u}>0$ on the whole $\bar{\Omega}$, and $u_{+}-\hat{u} \in \operatorname{int}\left(C_{+}\right)$. An analogous reasoning produces $\hat{u}-v_{-} \in \operatorname{int}\left(C_{+}\right)$. Hence,

$$
\begin{equation*}
\hat{u} \in \operatorname{int}_{C^{1}(\bar{\Omega})}\left(\left[v_{-}, u_{+}\right]\right) . \tag{4.8}
\end{equation*}
$$

A further nodal solution $\bar{u} \in C^{1}(\bar{\Omega})$ of (1.1) is easily obtained. Indeed, assertion $\left(\mathrm{j}_{2}\right)$ in Lemma 3.3 forces

$$
\begin{equation*}
C_{k}\left(\hat{\varphi}, u_{+}\right)=C_{k}\left(\hat{\varphi}, v_{-}\right)=\delta_{k, 0} \mathbb{Z}, \tag{4.9}
\end{equation*}
$$

while

$$
\begin{equation*}
C_{k}(\hat{\varphi}, 0)=\delta_{k, d_{m}} \mathbb{Z}, \quad \text { and } C_{k}(\hat{\varphi}, \infty)=\delta_{k, 0} \mathbb{Z} ; \tag{4.10}
\end{equation*}
$$

cf. the proof of Theorem 3.3. Now, if $K(\hat{\varphi})=\left\{0, u_{+}, v_{-}, \hat{u}\right\}$ then, combining (4.7), (4.9), and (4.10) with (2.1) we would immediately reach a contradiction. So, there exists $\bar{u} \in K(\hat{\varphi}) \backslash\left\{0, u_{+}, v_{-}, \hat{u}\right\}$. As already shown for $\hat{u}$, one has $\left.\bar{u} \in C^{1} \bar{\Omega}\right), \bar{u}$ nodal, besides

$$
\begin{equation*}
\bar{u} \in \operatorname{int}_{C^{1}(\bar{\Omega})}\left(\left[v_{-}, u_{+}\right]\right) . \tag{4.11}
\end{equation*}
$$

Finally, adapting the flow invariance arguments made in [8] we get a nodal solution

$$
\begin{equation*}
\tilde{u} \in C^{1}(\bar{\Omega}) \backslash \operatorname{int}_{C^{1}(\bar{\Omega})}\left(\left[v_{-}, u_{+}\right]\right) \tag{4.12}
\end{equation*}
$$

to problem (1.1). From (4.8), (4.11), and (4.12) it evidently follows $\tilde{u} \notin\{\hat{u}, \bar{u}\}$.
Remark 4.1. Let us note that $\tilde{u} \notin K(\hat{\varphi})$, otherwise, thanks to Lemma 3.3 and the trick at the beginning of the above proof, $\tilde{u} \in \operatorname{int}_{C^{1}(\bar{\Omega})}\left(\left[v_{-}, u_{+}\right]\right)$, against (4.12).

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